

Default Risk of a Time-homogeneous Diffusion Model^[1]

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2. A main result
3. On the auxiliary quantity $A(a, b, c)$
4. On the differentiability of $p(\cdot)$ at b
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- The traditional study of default risk usually defines default as the event that the firm value goes below a **too low level a** . In particular, $a = 0$ in classical ruin theory.
- In real world, however, the firm is also regarded as defaulted if its value constantly stays below a **moderately low level b** for a **certain time period c** , even though the hitting to level b does not immediately lead to illiquidity of the firm.
- Economic justifications for such a consideration are the **US bankruptcy codes** Chapter 7 (Liquidation) and Chapter 11 (Reorganization).

Default subject to two thresholds and one grace period

- Suppose that the value process is modeled by $X = \{X_t, t \geq 0\}$ with $X_0 = x_0$.
- For a real number x , denote by T_x the first hitting time to level x .
- Let $a < b$ and $c > 0$ be three exogenously determined constants, with a interpreted as the **liquidity threshold**, b as the **reorganization threshold** and c as the **grace period**.
- Let $\tau_b(c)$ be the first time when the process X has constantly stayed below level b for c units of time.
- See the graphs for \underline{T}_a and $\underline{\tau}_b(c)$.
- The **default probability** is defined by

$$q(x_0) = q(x_0; a, b, c) = \mathbf{P}^{x_0} \{T_a \wedge \tau_b(c) < \infty\}. \quad (1.1)$$

- It is sometimes more convenient to start with the **non-default probability** $p(x_0) = 1 - q(x_0)$.

Some immediate remarks

Letting $a \downarrow -\infty$ and $b = 0$ in (1.1) yields

$$q(x_0; -\infty, 0, c) = \mathbb{P}^{x_0} \{ \tau_0(c) < \infty \},$$

which is recognized as **Parisian ruin probability**.

Clearly, $q(x_0; a, b, c)$ is decreasing in c . Letting $c \downarrow 0$ in (1.1) yields

$$q(x_0; a, b, +0) = \mathbb{P}^{x_0} \{ T_b < \infty \},$$

while letting $c \uparrow \infty$ yields

$$q(x_0; a, b, \infty) = \mathbb{P}^{x_0} \{ T_a < \infty \}.$$

Hence, the grace period c serves as a bridge connecting the two **traditional default probabilities**:

$$\mathbb{P}^{x_0} \{ T_a < \infty \} \leq q(x_0; a, b, c) \leq \mathbb{P}^{x_0} \{ T_b < \infty \}$$

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The firm value process

Suppose that the firm value is modeled by a **time-homogeneous diffusion process** $X = \{X_t, t \geq 0\}$, with dynamics

$$dX_t = \mu(X_t)dt + \sigma(X_t)dW_t. \quad (2.1)$$

In (2.1):

- $X_0 = x_0$ is the initial wealth,
- $\{W_t, t \geq 0\}$ is a standard Brownian motion (Wiener process),
- and $\mu(\cdot)$ and $\sigma(\cdot) > 0$ are two measurable functions satisfying usual conditions of the existence and uniqueness theorem for the stochastic differential equation (2.1).

Denote by $\{\mathcal{F}_t, t \geq 0\}$ the natural filtration generated by $\{W_t, t \geq 0\}$.

The two-sided exit problem

Define

$$G(x) = \exp \left\{ - \int^x \frac{2\mu(y)}{\sigma^2(y)} dy \right\}, \quad S(x) = \int^x G(y) dy.$$

The function $S(\cdot)$ is referred to as the **scale function** of X . To avoid triviality, we assume that $S(\infty) < \infty$.

It is well known that, for $u < x < v$,

$$P^x \{ T_u < T_v \} = \frac{\int_x^v G(y) dy}{\int_u^v G(y) dy}, \quad P^x \{ T_u > T_v \} = \frac{\int_u^x G(y) dy}{\int_u^v G(y) dy}. \quad (2.2)$$

Letting $v = \infty$ in second relation in (2.2) yields

$$P^x \{ T_u = \infty \} = \frac{\int_u^x G(y) dy}{\int_u^\infty G(y) dy}. \quad (2.3)$$

The main result

Introduce an **auxiliary quantity**

$$A(a, b, c) = \lim_{\varepsilon \downarrow 0} \frac{\mathbb{P}^{b-\varepsilon} \{ T_b > T_a \wedge c \}}{\varepsilon}. \quad (2.4)$$

It will be proved later that $A(a, b, c)$ exists, is finite and equals the boundary derivative of the solution of a PDE. Hence, its value can be easily determined numerically.

Theorem 2.1 For $a < b \leq x_0$ and $c > 0$, the default probability satisfies

$$q(x_0) = \frac{A(a, b, c)}{A(a, b, c) \int_b^\infty G(y) dy + G(b)} \int_{x_0}^\infty G(y) dy. \quad (2.5)$$

Proof of Theorem 2.1

For $x \geq b$, by the strong Markov property,

$$\begin{aligned} p(x) &= \mathbb{P}^x \{ T_a = \infty, \tau_b(c) = \infty \} \\ &= \mathbb{P}^x \{ T_b = \infty \} + \mathbb{P}^x \{ T_a = \infty, \tau_b(c) = \infty, T_b < \infty \} \\ &= \mathbb{P}^x \{ T_b = \infty \} + \mathbb{E}^x [\mathbb{P}^x \{ T_a = \infty, \tau_b(c) = \infty, T_b < \infty | \mathcal{F}_{T_b} \}] \\ &= \mathbb{P}^x \{ T_b = \infty \} + \mathbb{P}^x \{ T_b < \infty \} p(b). \end{aligned} \tag{2.6}$$

It follows that

$$\begin{aligned} p'_+(b) &= \lim_{\varepsilon \downarrow 0} \frac{p(b + \varepsilon) - p(b)}{\varepsilon} \\ &= q(b) \lim_{\varepsilon \downarrow 0} \frac{\mathbb{P}^{b+\varepsilon} \{ T_b = \infty \}}{\varepsilon} \\ &= q(b) \frac{G(b)}{\int_b^\infty G(y) dy}, \end{aligned} \tag{2.7}$$

where in the last step we used (2.3).

Proof of Theorem 2.1 (Cont.)

Similarly, for $x \in (a, b)$ we have

$$p(x) = \mathbb{P}^x \{T_a = \infty, \tau_b(c) = \infty\} = \mathbb{P}^x \{T_b \leq T_a \wedge c\} p(b).$$

By Corollary 3.1 below, the limit $A(a, b, c)$ in (2.4) exists and is finite.

It follows that

$$\begin{aligned} p'_-(b) &= \lim_{\varepsilon \downarrow 0} \frac{p(b) - p(b - \varepsilon)}{\varepsilon} \\ &= p(b) \lim_{\varepsilon \downarrow 0} \frac{\mathbb{P}^{b-\varepsilon} \{T_b > T_a \wedge c\}}{\varepsilon} \\ &= p(b) A(a, b, c). \end{aligned} \tag{2.8}$$

Proof of Theorem 2.1 (Cont.)

By Theorem 4.1 below, the function $p(\cdot)$ is differentiable at b . Thus, the conjunction of (2.7) and (2.8) gives

$$p(b) = \frac{G(b)}{A(a, b, c) \int_b^\infty G(y) dy + G(b)}. \quad (2.9)$$

Substituting (2.9) into (2.6) and using (2.3), we obtain

$$p(x) = \frac{\int_b^x G(y) dy}{\int_b^\infty G(y) dy} + \frac{\int_x^\infty G(y) dy}{\int_b^\infty G(y) dy} \times \frac{G(b)}{A(a, b, c) \int_b^\infty G(y) dy + G(b)}.$$

Thus, relation (2.5) follows from $q(x) = 1 - p(x)$.

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Consider the modified two-sided exit probability function

$$\phi(x, t; a, b) = P^x \{T_b \leq T_a \wedge t\}, \quad a < x < b, t \geq 0.$$

The following theorem establishes a PDE for this function:

Theorem 3.1 Suppose $h(x, t)$ solves

$$h_t(x, t) = \mu(x)h_x(x, t) + \frac{1}{2}\sigma^2(x)h_{xx}(x, t), \quad a < x < b, t > 0,$$

with the boundary conditions $h(b, t) = 1$ and $h(a, t) = 0$ for $t \geq 0$ while $h(x, 0) = 0$ for $a < x < b$. Then

$$h(x, t) = \phi(x, t; a, b), \quad a \leq x \leq b, t \geq 0.$$

Proof of Theorem 3.1

Applying Itô's formula and noticing that $h_s(X_s, t - s) = -h_t(X_s, t - s)$,

$$\begin{aligned} & dh(X_s, t - s) \\ &= -h_t(X_s, t - s)ds + h_x(X_s, t - s)dX_s + \frac{1}{2}h_{xx}(X_s, t - s)\sigma^2(X_s)ds \\ &= h_x(X_s, t - s)\sigma(X_s)dW_s. \end{aligned}$$

For a stopping time τ , we have

$$h(X_{\tau \wedge t}, t - \tau \wedge t) = h(x, t) + \int_0^{\tau \wedge t} h_x(X_s, t - s)\sigma(X_s)dW_s.$$

This implies that $h(x, t) = \mathbf{E}^x [h(X_{\tau \wedge t}, t - \tau \wedge t)]$.

Let $\tau = T_a \wedge T_b$. By the boundary conditions, we have

$$\begin{aligned} h(x, t) &= \mathbf{E}^x [h(X_\tau, t - \tau)\mathbf{1}_{\{\tau \leq t\}}] + \mathbf{E}^x [h(X_t, 0)\mathbf{1}_{\{\tau > t\}}] \\ &= \mathbf{E}^x [h(a, t - T_a)\mathbf{1}_{\{T_a = \tau \leq t\}}] + \mathbf{E}^x [h(b, t - T_b)\mathbf{1}_{\{T_b = \tau \leq t\}}] + 0 \\ &= 0 + \mathbf{P}^x \{T_b \leq T_a \wedge t\} + 0. \end{aligned}$$

Existence and finiteness of $D(a,b,c)$

By the well-known regularity theory of PDE (see, e.g. Theorem 4.22 of Lieberman (1996)), we immediately have the following:

Corollary 3.1 It holds for every fixed $t > 0$ that $\phi_x(x, t; a, b)|_{x=b}$ is finite and continuous with respect to t . In particular,

$$\begin{aligned} A(a, b, c) &= \lim_{\varepsilon \downarrow 0} \frac{\mathbb{P}^{b-\varepsilon} \{T_b > T_a \wedge c\}}{\varepsilon} \\ &= \lim_{\varepsilon \downarrow 0} \frac{1 - \phi(b - \varepsilon, c; a, b)}{\varepsilon} \\ &= \phi_x(x, c; a, b)|_{x=b} \end{aligned}$$

exists and is finite.

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Laplace transforms of the two-sided exit times

Suppose that $g_1(\cdot; r)$ and $g_2(\cdot; r)$ are two independent positive solutions of

$$\frac{1}{2}\sigma^2(x)g''(x) + \mu(x)g'(x) = rg(x), \quad r \geq 0,$$

with $g_1(\cdot; r)$ decreasing and $g_2(\cdot; r)$ increasing. Define

$$f(y, z; r) = g_1(y; r)g_2(z; r) - g_1(z; r)g_2(y; r).$$

The Laplace transforms of the two-sided exit times for the diffusion process X were first solved by Darling and Siegert (1953):

Lemma 4.1 For $u \leq x \leq v$ and $r \geq 0$, we have

$$\mathbb{E}^x \left[e^{-rT_u}; T_u < T_v \right] = \frac{f(x, v; r)}{f(u, v; r)}, \quad \mathbb{E}^x \left[e^{-rT_v}; T_v < T_u \right] = \frac{f(u, x; r)}{f(u, v; r)}.$$

On the minimum of two exit times

Lemma 4.2 For all x ,

$$\lim_{\varepsilon \downarrow 0} \frac{\mathbb{E}^x [T_{x+\varepsilon} \wedge T_{x-\varepsilon}]}{\varepsilon^2} = C(x),$$

where $C(x) > 0$ is a finite constant with an explicit expression.

Proof. By Lemma 4.1,

$$\begin{aligned} \mathbb{E}^x [T_{x-\varepsilon}; T_{x-\varepsilon} < T_{x+\varepsilon}] &= - \left. \frac{\partial}{\partial r} \mathbb{E}^x \left[e^{-rT_{x-\varepsilon}}; T_{x-\varepsilon} < T_{x+\varepsilon} \right] \right|_{r=0} \\ &= - \left. \frac{\partial}{\partial r} \frac{f(x, x + \varepsilon; r)}{f(x - \varepsilon, x + \varepsilon; r)} \right|_{r=0} = \dots \end{aligned}$$

A similar relation for $\mathbb{E}^x [T_{x+\varepsilon}; T_{x+\varepsilon} < T_{x-\varepsilon}]$ holds. Therefore,

$$\begin{aligned} \mathbb{E}^x [T_{x+\varepsilon} \wedge T_{x-\varepsilon}] &= \mathbb{E}^x [T_{x+\varepsilon}; T_{x+\varepsilon} < T_{x-\varepsilon}] + \mathbb{E}^x [T_{x-\varepsilon}; T_{x-\varepsilon} < T_{x+\varepsilon}] \\ &= \text{Taylor's expansion} = \dots \end{aligned}$$

A lemma

Lemma 4.3 $\mathbb{P}^b \{ \tau_b(c) < T_{b-\varepsilon} \} = O(\varepsilon^2)$ as $\varepsilon \downarrow 0$.

Proof. For X starting with b , denote the consecutive down-crossing and up-crossing times of levels b and $b + \varepsilon$ by

$$0 = \gamma_0^- < \gamma_0^+ < \cdots < \gamma_{i-1}^- < \gamma_{i-1}^+ < \gamma_i^+ < \cdots .$$

Then, by the strong Markov property of X and Lemma 4.2,

$$\begin{aligned} \mathbb{P}^b \{ \tau_b(c) < T_{b-\varepsilon} \} &= \sum_{i=0}^{\infty} \mathbb{P}^b \{ \gamma_i^- < \tau_b(c) < \gamma_i^+ \wedge T_{b-\varepsilon} \} \\ &= \sum_{i=0}^{\infty} \mathbb{E}^b \left[\mathbb{E}^b \left[\mathbf{1}_{\{ \gamma_i^- < \tau_b(c) < \gamma_i^+ \wedge T_{b-\varepsilon} \}} \middle| \mathcal{F}_{\gamma_i^-} \right] \right] \\ &= \mathbb{P}^b \{ \tau_b(c) < T_{b+\varepsilon} \wedge T_{b-\varepsilon} \} \sum_{i=0}^{\infty} \mathbb{P}^b \{ \gamma_i^- < T_{b-\varepsilon} \} \\ &\leq \frac{1}{c} \mathbb{E}^b [T_{b+\varepsilon} \wedge T_{b-\varepsilon}] \sum_{i=0}^{\infty} \left(\mathbb{P}^b \{ T_{b+\varepsilon} < T_{b-\varepsilon} \} \right)^i \\ &= O(\varepsilon^2). \end{aligned}$$

Differentiability of $p(x)$ is at b

Theorem 4.1 The function $p(\cdot)$ is differentiable at b .

Proof. By (2.7), it suffices to show that

$$\lim_{\varepsilon \downarrow 0} \frac{p(b) - p(b - \varepsilon)}{\varepsilon} = q(b) \frac{G(b)}{\int_b^\infty G(y) dy}. \quad (4.1)$$

Proof of Theorem 4.1 (Cont.)

For small $\varepsilon > 0$, we split $p(b)$ into two parts as

$$P^b \{ \tau_b(c) = \infty, T_{b-\varepsilon} = \infty \} + P^b \{ T_a = \infty, \tau_b(c) = \infty, T_{b-\varepsilon} < \infty \}.$$

By Lemma 4.3, the first term equals

$$P^b \{ T_{b-\varepsilon} = \infty \} - P^b \{ \tau_b(c) < \infty, T_{b-\varepsilon} = \infty \} = P^b \{ T_{b-\varepsilon} = \infty \} + O(\varepsilon^2).$$

With some efforts, we can prove that the second term satisfies

$$P^b \{ T_{b-\varepsilon} < \infty \} p(b - \varepsilon) + o(\varepsilon)$$

It follows that

$$p(b) = P^b \{ T_{b-\varepsilon} = \infty \} + P^b \{ T_{b-\varepsilon} < \infty \} p(b - \varepsilon) + o(\varepsilon),$$

which easily implies (4.1).

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Implicit finite difference method

Recall formula (2.5) for the default probability $q(x_0)$:

$$q(x_0) = \frac{A(a, b, c)}{A(a, b, c) \int_b^\infty G(y)dy + G(b)} \int_{x_0}^\infty G(y)dy. \quad (2.5)$$

The only implicit part is the quantity $A(a, b, c)$.

Theorem 3.1 and Corollary 3.1 enable us to compute $A(a, b, c)$ numerically via a PDE. We use the second-order **implicit finite difference method** to solve $A(a, b, c)$ for both a Brownian motion and a geometric Brownian motion.

Brownian motion case

Assume $dX_t = \mu dt + \sigma dW_t$, with $x_0 \geq b$, $\mu, \sigma > 0$. Relation (2.5) reduces to (with $\rho = 2\mu/\sigma^2$)

$$q(x_0) = \frac{A(a, b, c)}{A(a, b, c) + \rho} e^{-\rho(x_0 - b)}.$$

Set parameters to $\mu = 0.1$, $\sigma = 0.25$, $a = 0.1$, $b = 0.2$ and $c = 1$. Then

mesh	$A(0.1, 0.2, 1)$	$q(x_0)$	elapsed time (s)
0.005	8.5534038	$1.3801420 \times e^{-3.2x_0}$	0.505305
0.001	8.4987776	$1.3777311 \times e^{-3.2x_0}$	4.119148
0.0005	8.4919795	$1.3774294 \times e^{-3.2x_0}$	44.813284
0.00025	8.4885830	$1.3772786 \times e^{-3.2x_0}$	342.789096

In the meantime,

$$P^{x_0} \{T_a < \infty\} = 1.3771278 \times e^{-3.2x_0},$$

$$P^{x_0} \{T_b < \infty\} = 1.8964809 \times e^{-3.2x_0},$$

$$P^{x_0} \{\tau_b(c) < \infty\} = 0.6932042 \times e^{-3.2x_0}.$$

Geometric Brownian motion case

Assume $dX_t = \mu X_t dt + \sigma X_t dW_t$, with $x_0 \geq b$ and $\rho = 2\mu/\sigma^2 > 1$.
Relation (2.5) reduces to

$$q(x_0) = \frac{A(a, b, c)}{A(a, b, c) + \rho - 1} \left(\frac{b}{x_0}\right)^{\rho-1}.$$

Set parameters to $\mu = 0.1$, $\sigma = 0.25$, $a = 0.1$, $b = 0.2$ and $c = 1$. Then

mesh	$A(0.1, 0.2, 1)$	$q(x_0)$	elapsed time (s)
0.005	11.495846	$0.024334255 \times x_0^{-2.2}$	0.460361
0.001	11.145638	$0.024212051 \times x_0^{-2.2}$	4.881175
0.0005	11.101517	$0.024196199 \times x_0^{-2.2}$	37.679470
0.00025	11.079429	$0.024188223 \times x_0^{-2.2}$	287.741042

In the meantime,

$$P^{x_0} \{T_a < \infty\} = 0.0063095734 \times x_0^{-2.2},$$

$$P^{x_0} \{T_b < \infty\} = 0.028991187 \times x_0^{-2.2},$$

$$P^{x_0} \{\tau_b(c) < \infty\} = 0.014533261 \times x_0^{-2.2}.$$

Potential future works

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- Incorporate (heavy-tailed) jumps into the modeling;
- Consider the more practical finite-time case.
- ...

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Thank You Very Much!!!